

**ELASTIK YOPISHQOQ SUYUQLIKLARNING STASIONAR BO’LMAGAN  
OQIMLARI HARAKATINI O’RGANISHDA SHULMAN-XUSIDNING MODIFIKASIYA  
MODELINI QO’LLASH.**

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**Annotasiya.** Kanal va quvurlarda elastik yopishqoq suyuqliklarning statsionar bo‘lmagan va pulsatsiyalanuvchi oqimlarini o’rganishda ularning murakkab reologik modellarini qo’llash orqali aniq amaliy masalalarni yechish, analitik usullarni ishlab chiqish ishlari olib borilgan va ijobiy natijalarga erishilgan. Ushbu maqolada biz kanallar va quvurlardagi elastik yopishqoq suyuqliklarning stosionar bo’lmagan oqimlari harakatlarini o’rganish uchun Shulman-Xusid modelining modifikasiyasidan foydalanishni taklif qilamiz va bu model Nyuton va Maksvell modellarining umumiy holi ekanligini ko’rsatamiz.

**ПРИМЕНЕНИЕ МОДИФИКАЦИОННОЙ МОДЕЛИ ШУЛЬМАНА-ХУСИДА  
ПРИ ИССЛЕДОВАНИИ ДВИЖЕНИЯ НЕСТАЦИОНАРНЫХ ПОТОКОВ  
УПРУГИХ ВЯЗКИХ ЖИДКОСТЕЙ.**

**Аннотация.** При исследовании нестационарных и пульсирующих течений упругих вязких жидкостей в каналах и трубах с применением их сложных реологических моделей проведены работы по решению конкретных практических задач, разработке аналитических методов и достигнуты положительные результаты. В данной статье мы предлагаем использовать модификацию модели Шульмана-Хусида для исследования поведения нестационарных течений упругих вязких жидкостей в каналах и трубах и показываем, что эта модель является общим случаем моделей Ньютона и Максвелла.

**APPLICATION OF SHULMAN-KHUSID MODIFICATION MODEL IN STUDYING THE  
MOVEMENT OF NON-STATIONARY FLOWS OF ELASTIC VISCOUS FLUIDS.**

**Abstract.** In the study of non-stationary and pulsating flows of elastic viscous fluids in channels and pipes, by applying their complex rheological models, work was carried out to solve specific practical problems, to develop analytical methods, and positive results were achieved. In this article, we propose to use a modification of the Shulman-

Husid model to study the behavior of non-stationary flows of elastic viscous fluids in channels and pipes, and we show that this model is a general case of the Newtonian and Maxwellian models.

**Tayanch so'zlar.** Elastik yopishqoq suyuqlik, kanal, quvur, stosionar bo'limgan oqim, model, reloksasiya, spektr, modifikatsiya.

**Ключевые слова:** Упругая вязкая жидкость, канал, труба, нестационарное течение, модель, релоксация, спектр, модификация.

**Key words:** Elastic viscous fluid, channel, pipe, non-stationary flow, model, reallocation, spectrum, modification.

Kirish (Introduction) Mavzuga oid adabiyotlarning tahlili (Literature review). Elastik yopishqoq suyuqliklarning matematik modeli asosan suyuqlik elastikligining o'zgarishini hisobga olishga asoslangan. Eng umumiyl holda muhit holatning reologik tenglamalari Koshi va Finger chekli deformatsiyalari tenzor bog'lanishi bilan ifodalangan kuchlanishlar tenzori ko'rinishida berilgan [14, 15]. Keyin elastik yopishqoq suyuqliklarning holatlari chiziqsiz reologik tenglamalari Maksvell modeli ko'rinishida umumlashtirilgan [1, 6, 7, 17, 20]. U yerda relaksasion va integral tipdagi holatlar chiziqsiz elastik yopishqoq reologik tenglamalarning matematik ekvivalentligi shartlari asosida aniqlanadi.

Yassi va xalqasimon doiraviy kesimli uzun kanallarda elastik yopishqoq suyuqlik oqimi Z.P.Shulman, B.M. Xusidlar tomonidan tadqiq etilgan [17].

Bosim gradiyentining vaqt bo'yicha o'zgarishi natijasidagi yopishqoq Nyuton suyuqliklarining laminar nostasionar va stasionar oqimlar masalalari [2, 3, 4, 8] tadqiqot ishlarida qaralgan bo'lsa, turbulent rejimdagi oqimlar [13, 16] ishlarda tadqiq qilingan. Nyuton suyuqligi harakatini ko'pfazoli muhit sifatida qarash, bu sohadagi nazariy tadqiqot ishlarining tajriba natijalariga yaqinlashtirish imkonini beradi, bunday tadqiqotlar [5, 13] ishda mukammal va amaliyotga tadbiqlari asosida keltirilgan.

Ammo nonyuton suyuqliklar harakatini o'rganishda ko'plab tadqiqotlar olib borilganiga qaramay, hozirgi kunda Shulman-Xusid modeli yordamida masalalar yetarli darajada tadqiq qilinmagan va amaliyotga tadbiq etilmagan. Jumladan, elastik yopishqoq suyuqliklarning oqimida, bosim gradiyentining keskin o'sishi va keskin kamayib nolga teng bo'lishidagi nostasionar oqimlarida sodir bo'ladigan gidrodinamik o'zgarishlar batafsil o'rganilmagan [10, 14, 17, 18, 19, 20].

Tadqiqot metodologiyasi (Research Methodology). Ma'lumki, ko'pchilik hollarda elastik yopishqoq suyuqlik oqimi uchun Maksvellning bir o'lchovli fazodagi klassik modeli elastik yopishqoq suyuqliklar oqimi uchun qo'lanilib kelingan [1, 8, 9]:

$$\tau + \lambda \dot{\tau} = 2\eta D, \quad (1)$$

bu yerda  $\tau$  – kuchlanishlar tenzori;  $\dot{\tau}$  – kuchlanishlar tenzorining vaqt bo'yicha hosilasi;  $\eta$  – suyuqlikning dinamik yopishqoqlik koeffisenti;

$\lambda = \frac{\eta}{E}$  – relaksatsiya koeffisiyenti;  $E$  – yopishqoq suyuqlikning elastiklik moduli;

$D = \frac{1}{2} \frac{\partial u}{\partial y}$  – muhitning siljish deformatsiyasi tezligi.

Biz Maksvell modeli ko‘rinishidagi barcha elastik yopishqoq suyuqliklarning reologik modellarini topologik model shaklida umumlashtiruvchi Shulman-Xusidning modelida keltiramiz:

$$\begin{aligned} T &= \sum_{k=1}^{\infty} \left( 1 + \frac{\varepsilon}{2} \right)' T_k^{(1)} + \frac{\varepsilon}{2} T_k^{(2)}, \quad T_k^{(1)} + \frac{\nabla^{(1)}}{\lambda_k} T_k^{(1)} = 2 p_k \mathbf{D}, \\ T_k^{(2)} &+ \frac{\nabla^{(2)}}{\lambda_k} T_k^{(2)} = -2 p_k \mathbf{D}, \quad \frac{D p_k}{D t} + \frac{g_k}{\lambda_k} p_k = \frac{\eta_k}{\lambda_k^2} f_k. \end{aligned} \quad (2)$$

Bu yerda yuqori konvektiv hosila ushbu ifoda orqali

$$T_k^{(1)} = \frac{DT_k^{(1)}}{Dt} - T_k^{(1)} \nabla V^T - \nabla V \cdot T_k^{(1)},$$

pastki konvektiv hosila esa quyidagi ifoda orqali aniqlanadi:

$$T_k^{(2)} = \frac{DT_k^{(2)}}{Dt} + T_k^{(2)} \nabla V + \nabla V^T \cdot T_k^{(2)}.$$

Yaumann hosilasi esa ushbu ko‘rinishda berilgan:  $\frac{DA}{Dt} = \frac{\partial A}{\partial t} + V \nabla A + WA - AW$ ,

$$\text{bu yerda } \nabla V = \mathbf{D} + W, \quad \mathbf{D} = \frac{1}{2} (\nabla V^T + \nabla V), \quad W = \frac{1}{2} (\nabla V - \nabla V^T) \quad \text{bo‘lib},$$

$\mathbf{D}$  – deformatsiya tezligi tenzori;  $\nabla V$  – deformatsiya tezligi gradiyenti;  $\nabla V^T$  – transponirlangan deformatsiya tezligi gradiyenti; noldanfarqli ikkinchi normal kuchlanishlar ayirmasini ifodalovchi  $\varepsilon$  parametr sifatida kiritilgan bo‘lib, u quyidagi

formula orqali aniqlanadi  $\frac{\varepsilon}{2} = \frac{\Psi_2}{\Psi_1}$ . Bu yerda  $\Psi_1 = (\sigma_{11} - \sigma_{22})/\gamma^2$ ,  $\Psi_2 = (\sigma_{22} - \sigma_{33})/\gamma^2$

funksiyalar, mos ravishda, birinchi va ikkinchi, ikkinchi va uchinchi kuchlanishlar ayirmasi;  $p_k$  – tenglamaning o‘zidan aniqlanuvchi parametr;  $T_k^{(1)}$ ,  $T_k^{(2)}$  – har biri uchun

to‘qqizta elementdan tashkil topgan kuchlanish tenzorlari;  $T_k^{(1)}$ ,  $T_k^{(2)}$  – ixtiyoriy tanlangan koordinatalar sistemasida berilgan kuchlanish tenzorlari bo‘lib, ular uchun belgilangan hosilalar, mos ravishda, yuqori va pastki konvektiv hosilalarni bildiradi [18-19].

Yuqorida keltirilgan (2) umumiyoq ko‘rinishdagi elastik yopishqoq suyuqliklar modeliga juda ko‘p polimer suyuqliklar va boshqa elastik yopishqoq suyuqliklar modellari kiradi. Asosan, bu modellarning bir-biridan farqi, elastik yopishqoq suyuqliklar xususiyatlariga qarab (2) tenglamaga kiruvchi  $f_k(S_D(t'))$  va  $g_k(S_D(t'))$

funksiyalarning aniqlanishidadir. Xususan, kichik deformasiyalarda  $f_k = g_k = 1$  bo‘lib, bu holda elastik yopishqoq suyuqliklar modeli chiziqli ko‘rinishga keladi. Sonli hisob yuritishda  $\lambda_k$  va  $\eta_k$  kattaliklar, jumladan,  $\lambda_k = \frac{\lambda}{k^\alpha}$ ,  $\eta_k = \frac{\eta}{\xi(\alpha)k^\alpha}$  ko‘rinishida olinadi, bu yerda  $\eta$  – boshlang‘ich holatdagi Nyuton suyuqligining dinamik yopishqoqlik koeffisiyenti;  $\lambda$  – relaksasiya koeffisiyenti (vaqt);  $\alpha$  – relaksasiya vaqtini taqsimlanish spektrini xarakterlovchi son;  $\xi(\alpha)$  – Rimanning dzeta funksiyasi, u  $\xi(\alpha) = \sum_{k=1}^{\infty} \frac{1}{k^\alpha}$  ko‘rinishidagi ifoda orqali aniqlanadi.

Shuni aytish kerakki, yuqori konvektiv hosila, pastki konvektiv hosila va Yaumann hosilasi ixtiyoriy koordinatalar sistemasida olingan bo‘lib, ular ortogonal Dekart va silindrik koordinatalar sistemalarida oddiy hosilalarga aylanadi.

Maksvell tipidagi modellarda  $f_k = g_k = 1$  va  $p_k = \frac{\eta_k}{\lambda_k}$  ekanligidan, yuqorida keltirilgan Shulman-Xusid modeli to‘liq chiziqli ko‘rinishga keladi. Shunday ekan Shulman-Xusid modelini, ushbu ko‘rinishdagi modifikatsiyalangan model shakliga keltiramiz[10, 11]:

**Dekart koordinatalar sistemasi uchun ushbu ko‘rinishda:**

$$T = \sum_{k=1}^{\infty} T_k, \quad \lambda_k \frac{\partial T_k}{\partial t} + T_k = 2\eta_k \mathbf{D}, \\ \mathbf{D} = \frac{1}{2} \frac{\partial \mathbf{u}}{\partial \mathbf{y}}; \quad (3)$$

**silindrik koordinatalar sistemasi uchun esa ushbu ko‘rinishda keltiramiz:**

$$T = \sum_{k=1}^{\infty} T_k, \quad \lambda_k \frac{\partial T_k}{\partial t} + T_k = 2\eta_k \mathbf{D}, \\ \mathbf{D} = \frac{1}{2} \frac{\partial \mathbf{g}_x}{\partial r}. \quad (4)$$

Bu yerda keltirilgan modifikatsiyalangan Shulman-Xusid modelidan xususiy holda Nyuton, Maksvell modellari kelib chiqishini keltiramiz. Ryelaksasiya koeffisiyenti  $\lambda$  nolga intilgan holatida xususiy hol sifatida Nyuton modeli kelib chiqadi. Haqiqatdan ham, (3) va (4) tenglamalarda  $\lim_{\lambda \rightarrow 0} \lambda_k = 0$  desak, u holda Dekart koordinatalar sistemasi uchun

$$T = \eta \frac{\partial \mathbf{u}}{\partial \mathbf{y}} \quad (5)$$

va silindrik koordinatalar sistemasi uchun

$$T = \eta \frac{\partial \mathcal{G}_x}{\partial r} \quad (6)$$

ko‘rinishda ifodalash qiyin emas.

Modifikatsiyalangan Shulman-Xusid modelining xususiy holi sifatida, ryelaksasiya vaqtiga spektrini taqsimlanishini xarakterlovchi son  $\alpha$  cheksizlikka intilganda, Maksvell suyuqligi uchun reologik tenglama hosil bo‘ladi. Ya’ni (3) va (4) tenglamalarda  $\alpha \rightarrow \infty$  bo‘lganda, Dekart koordinatalar sistemasida

$$\tau + \lambda \dot{\tau} = \eta \frac{\partial u}{\partial y} \quad (7)$$

tenglamani va silindrik koordinatalar sistemasida

$$\tau + \lambda \dot{\tau} = \eta \frac{\partial \mathcal{G}_x}{\partial r} \quad (8)$$

tenglamani hosil qilamiz. Bu tenglamalar Maksvell tenglamalarini ifoda qiladi.

Tahlil va natijalar (Analysis and results). Nyuton modelining kamchiligi suyuqlikning elastiklik xususiyatini e’tiborga olmasligida, afsuski elastiklik xususiyatini o‘z ichiga oluvchi suyuqliklar texnika va texnologik jarayonlarda juda ko‘plab uchraydi [20].

Keltirilgan Maksvell modelining Shulman-Xusid modeliga nisbatan kamchiligi shundaki, bu yerda elastik yopishqoq suyuqlik bir jinsli suyuqlikdan iborat deb, faqat bitta relaksatsiya vaqtiga (koeffisiyenti) orqali aniqlanadi. Shulman-Xusid modelining modifikatsiyasida esa elastik yopishqoq suyuqlik bir jinsli emas deb qaraladi va relaksatsiya vaqtiga (koeffisiyenti) cheksiz ko‘p yaqinlashuvchi ketma-ketlik sifatida ifodalanib, bular ichidagi eng katta relaksatsiya koeffisiyenti tanlanadi, undan keyingi relaksatsiya koeffisiyentlari esa uning ma’lum bir qonuniyatiga bo‘ysinuvchi ulushi sifatida aniqlanadi. Shu boisdan ham Shulman-Xusid modelining modifikatsiyasi Maksvell modelininining umumlashmasi sifatida qaraladi [17-20].

Xulosa (Conclusion). Yuqorida taklif etilgan Shulman-Xusidning modifikatsiya modeli orqali yassi kanalda elastik yopishqoq suyuqlikning stasionar bo‘lmagan oqimi harakatini o‘rganish Maksvell modelida elastik yopishqoq suyuqlikning stasionar bo‘lmagan oqimi harakatini urganishdan bir qancha qulay va yaxshi natija berishini ko‘rish mumkin.

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